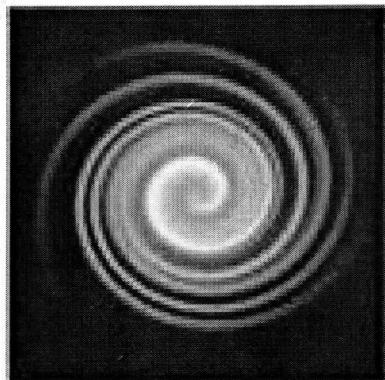


Application of a slender vortex filament code to the study of a four-vortex wake model

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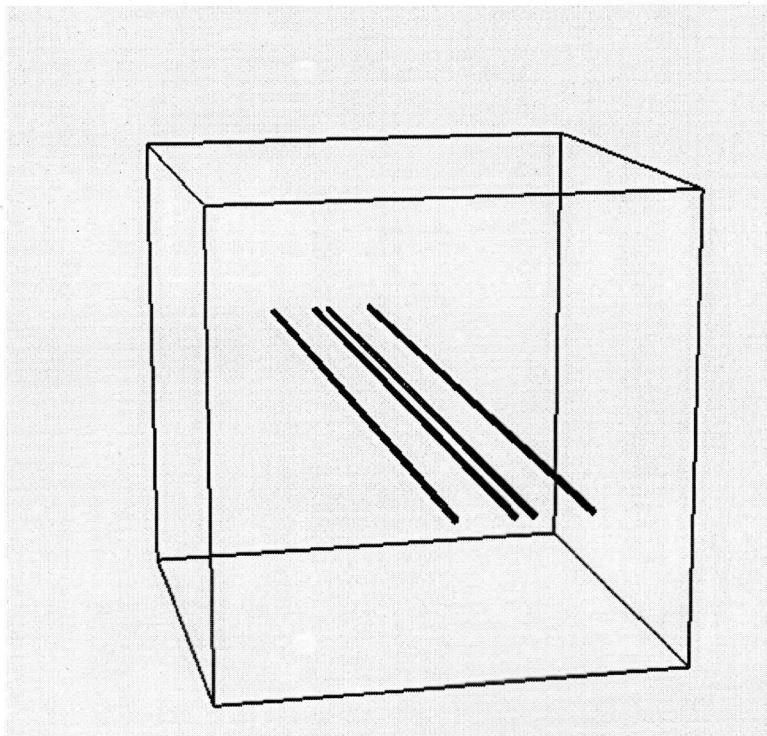
Fluid Mechanics Institute of Toulouse

C-Wake european project GRD1-1999-10332

Papers

1. Implementation and validation of a slender vortex filament code, *submitted to Int. J. Numer. Meth. Fluids.*
2. EZ-Vortex documentation: a slender vortex filament solver.
3. Effects of simple generic configuration on near to far field wake by using 3D vortex filament method, C-Wake *partner report UPS-PR 2.2.3-2.*

Introduction:

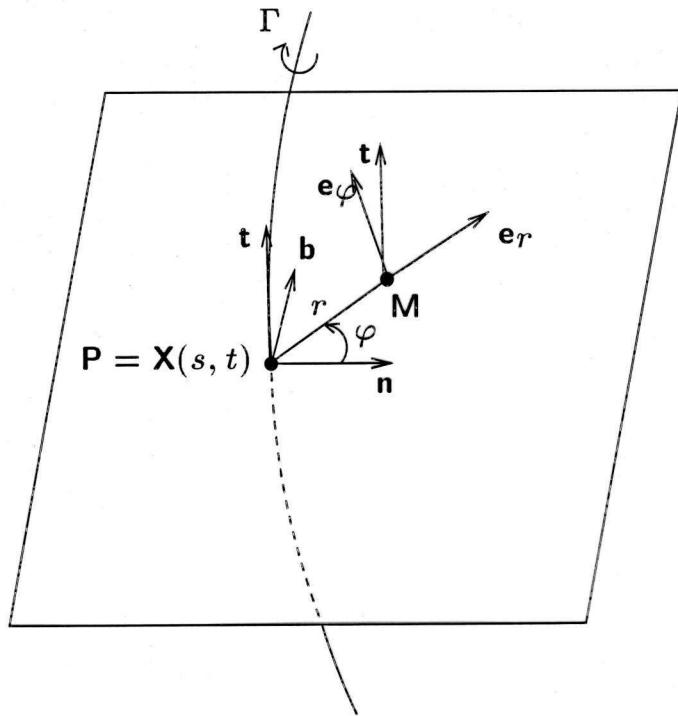


The four-vortex wake

- Rennich and Lele (AIAA 1998): simple vortex filament method
- Fabre and Jacquin (Phy. Fluids 2000): linear stability study
- Klein and Knio (J. Comput. 2000): new improved numerical schemes for slender vortex simulations (vortex ring)

⇒ implement the Klein and Knio scheme for open filaments (periodicity)
⇒ comparison with linear stability results

Equation of motion of slender filaments



The vortex filament.

Cut-off Equations

$$\frac{\partial \mathbf{X}}{\partial t} = \frac{1}{4\pi} \int_I \sigma(s', t) \frac{\mathbf{t}(s', t) \times (\mathbf{X}(s, t) - \mathbf{X}(s', t))}{|\mathbf{X}(s, t) - \mathbf{X}(s', t)|^3} ds'$$

where $I = [0, 2\pi[\setminus [s - s_c, s + s_c[$ s_c : cut-off length
 \Rightarrow Ad-hoc method. How to chose s_c ?
 \Rightarrow Callegari and Ting: Siam Applied Math. 78

$$s_c(s, t) = \epsilon \exp(1 - \ln 2 - C_v(t) - C_w(t)) / \sigma(s, t)$$

The equation of Callegari and Ting

$$\boxed{\partial \mathbf{X}/\partial t = \mathbf{A} + \frac{\Gamma K(s, t)}{4\pi} B \mathbf{b}(s, t),}$$

where

$$\mathbf{A}(s, t) = \frac{\Gamma}{4\pi} \int_{-\pi}^{+\pi} \sigma(s + s', t) \mathbf{N} ds',$$

$$\mathbf{N} = \frac{\mathbf{t}(s + s', t) \times (\mathbf{X}(s, t) - \mathbf{X}(s + s', t))}{|\mathbf{X}(s, t) - \mathbf{X}(s + s', t)|^3} - \frac{K(s, t) \mathbf{b}(s, t)}{2 |\lambda(s, s', t)|},$$

$$\sigma(s, t) = |\partial \mathbf{X}/\partial s|,$$

$$\lambda(s, s', t) = \int_s^{s+s'} \sigma(s^*, t) ds^*,$$

and

$$B = -\log \epsilon + \log(S) - 1 + C_v(t) + C_w(t)$$

K : local curvature

S : length of the filament

ϵ : small dimensionless thickness

Similar vortex core

$$\begin{aligned} C_v(t) &= [1 + \gamma - \ln 2]/2 - \ln(\bar{\delta}), \\ C_w(t) &= -2(S_0/S)^4(m_0/(\Gamma\bar{\delta}))^2, \end{aligned}$$

where

$$\begin{aligned} \bar{\delta}^2(t) &= \bar{\delta}_0^2 \left(\frac{S_0}{S(t)} \right) \left(1 + \frac{\bar{\delta}_{\bar{\nu}}^2}{\bar{\delta}_0^2} \right) \\ \bar{\delta}_{\bar{\nu}}^2 &= 4\bar{\nu} \int_0^t \frac{S(t')}{S_0} dt', \end{aligned}$$

γ = Euler number. Subscript 0 stands for initial.

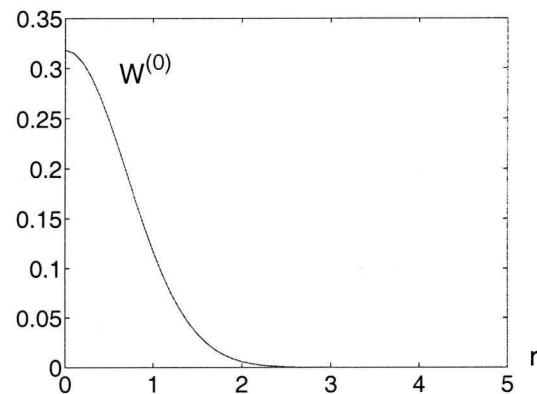
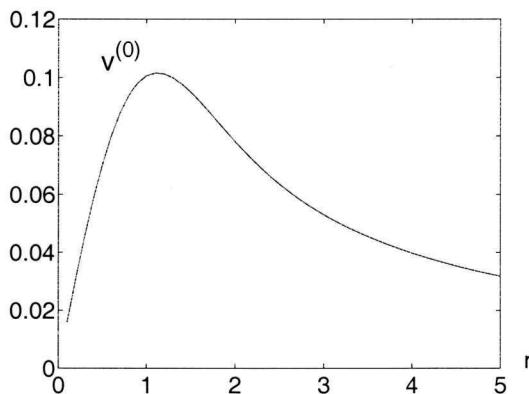
$\bar{\delta} = \delta/\epsilon$: the stretched radius

$\bar{\nu} = \nu/\epsilon^2$ is the stretched kinematic viscosity.

m_0 : the initial axial flux of the ring.

Circumferential $v^{(0)}$ and axial $w^{(0)}$ velocities:

$$v^{(0)} = \frac{\Gamma}{2\pi\bar{r}} \left[1 - e^{-\bar{r}/\bar{\delta}} \right], \quad w^{(0)} = \frac{m_0}{\pi\bar{\delta}^2} \left(\frac{S_0}{S} \right)^2 e^{-\bar{r}/\bar{\delta}},$$

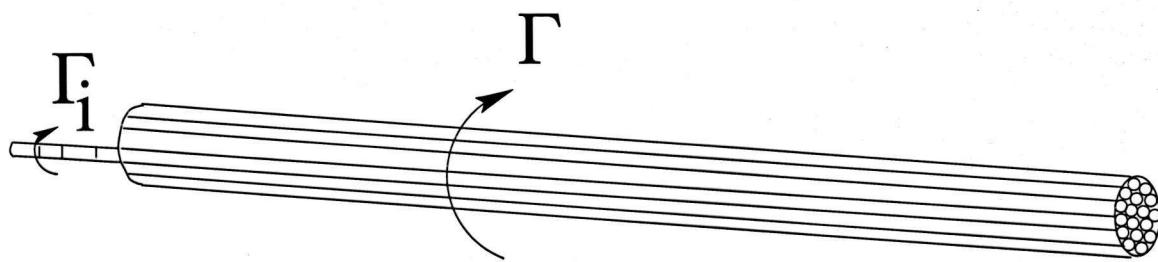


$\bar{r} = r/\epsilon$: the stretched radial distance to the filament,

Vortex Methods / Slender vortex

Methods

1) Vortex Blob or Vortex filament Methods

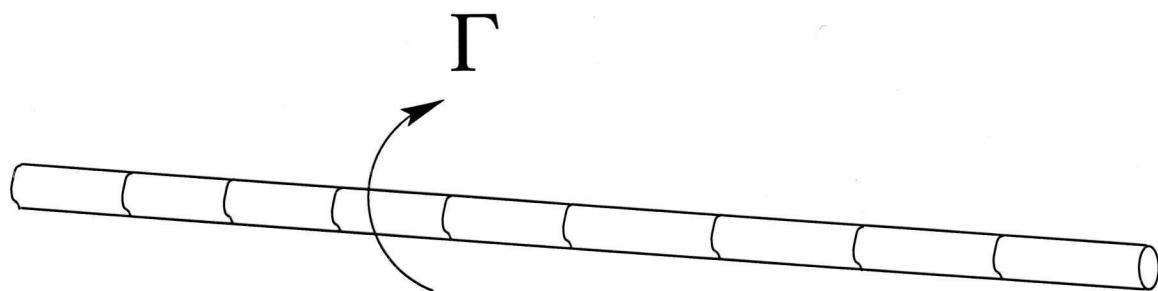


The computation with 1 filament per section is not correct (Klein and Knio JFM 95)

⇒ Great number of filaments per section to converge or

⇒ Theoretical correction of the method (Klein and Knio JFM 95/Klein and Knio 2000): new slender vortex solver

2) Slender Vortex filament solver:



+ Nbr filaments per section = 1

+ Not stiff with thickness:

boundary layer solved theoretically

⇒ Small thickness, no reconnection, no short wavelengths

M1 method of Knio and Klein

$$\boxed{\partial \mathbf{X}/\partial t = \mathbf{v}_{\sigma_1} + (\mathbf{v}_{\sigma_1} - \mathbf{v}_{\sigma_2}) \frac{\log(\sigma_1/\delta^{ttm})}{\log(\sigma_2/\sigma_1)}}$$

where

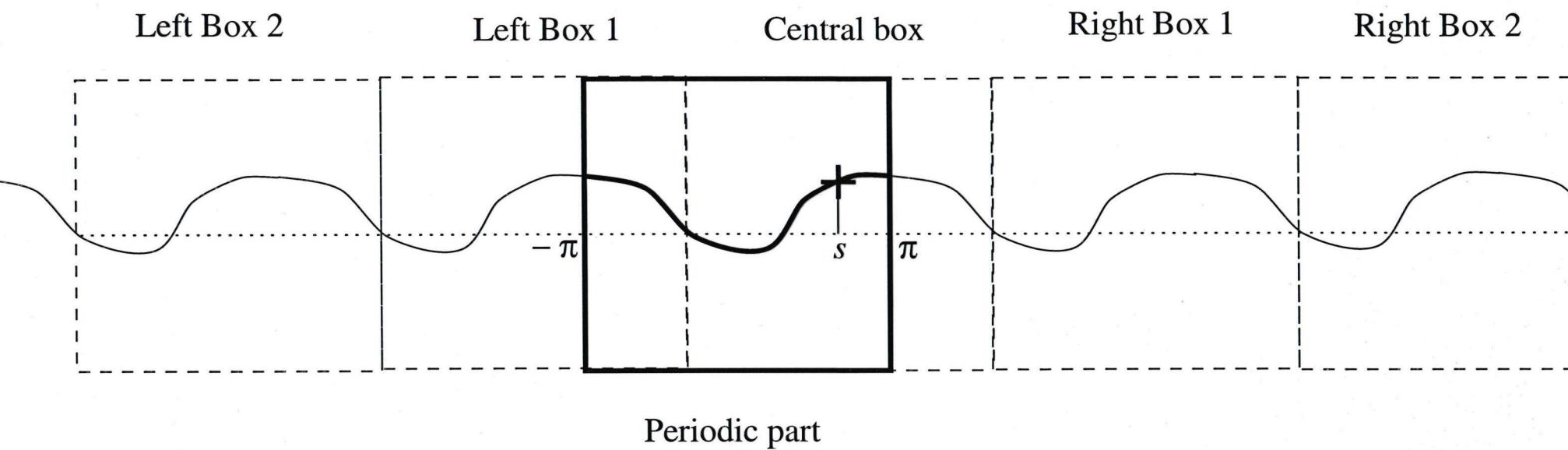
$$\begin{aligned}\mathbf{v}_x &= \frac{\Gamma}{4\pi} \int_{\mathcal{C}} \sigma(s', t) \mathbf{N}_x ds', \\ \mathbf{N}_x &= \frac{\mathbf{t}(s', t) \times (\mathbf{X}(s, t) - \mathbf{X}(s', t))}{\left[|\mathbf{X}(s, t) - \mathbf{X}(s', t)|^2 \right]^{3/2}} \kappa \left(\frac{|\mathbf{X}(s, t) - \mathbf{X}(s', t)|}{x} \right) \\ \kappa(r) &= \tanh(r^3)\end{aligned}$$

and

$$\begin{aligned}\sigma_1 &= 3\sigma_{max}, \\ \sigma_2 &= 2\sigma_{max}, \\ \sigma_{max} &= ds \max_{s \in [0, 2\pi]} \sigma(s, t), \\ \delta^{ttm} &= \epsilon \exp(C^{ttm} + 1 - C_v(t) - C_w(t)),\end{aligned}$$

With the choice of $\kappa(r) = \tanh(r^3)$, the C^{ttm} constant is $C^{ttm} = -0.4202$.

Implementation: open filament (periodic)



code EZ-vortex: - M1 méthod / CST equation / LIA
↳ fast, simple
well documented
easy to handle.

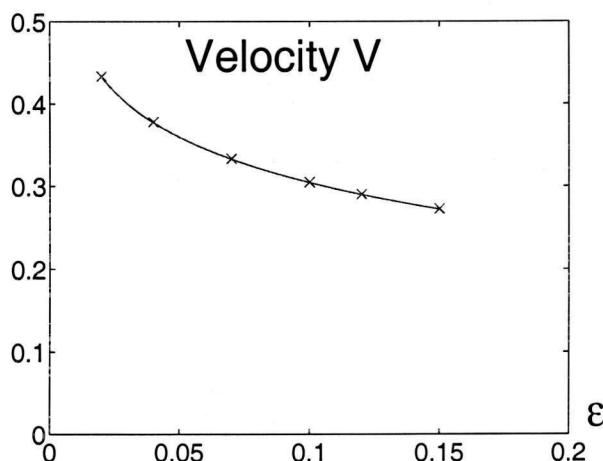
EZ-vortex¹: a Slender Vortex Filament solver (SVF)

- 1) Input: Initial position of the filaments
- 2) Solver: SVF solver for *closed* or *open* filaments
 - Equation:
 - The Callegari and Ting Equation (Implicit stepping)
 - The M1 method of Knio and Klein (Adams-Basforth or Implicit)
 - Spatial derivatives: Finit. Diff. or Spectral
 - Core: similar core or non-similar (Laguerre series)
 - Inviscid or viscous
- 3) Output:
 - Run-time drawing with OpenGL (SGI or linux)
 - File *history.dat* and mode 'Visualisation' of this file
 - Movie of the simulation

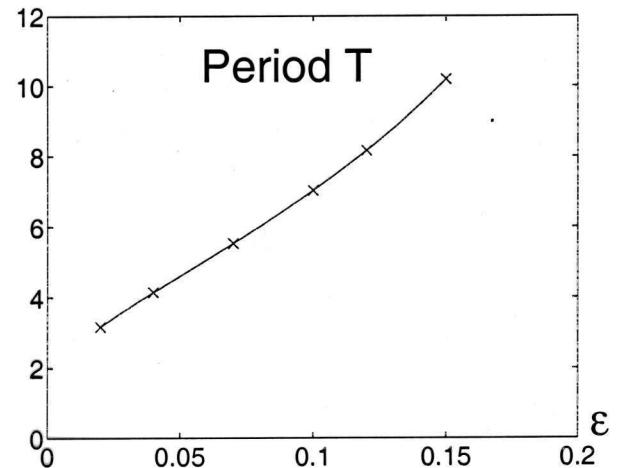
¹ D.Margerit, A.Giovannini, P. Brancher: EZ-vortex documentation: a Slender Vortex Filament solver

Validation

Linear stability analysis of a perturbed circular vortex ring (similar vortex and inviscid):

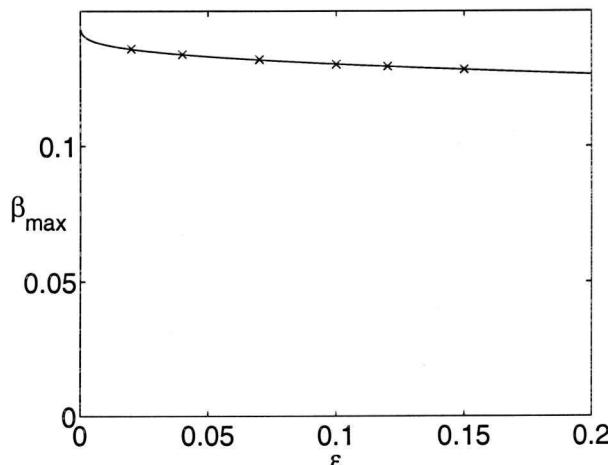


Velocity $V = f(\epsilon)$

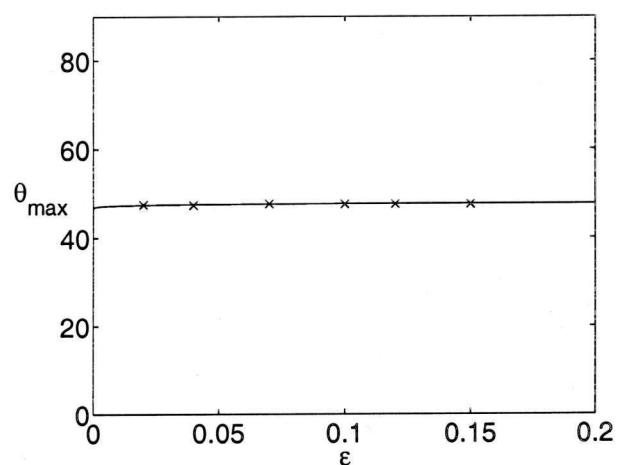


Period $T = f(\epsilon)$ for mode 3

Linear stability analysis of a pair of contra-rotating vortex filaments (similar vortex and inviscid):



Ampl. rate $\beta_{max} = f(\epsilon)$



Planar angle $\theta_{max} = f(\epsilon)$

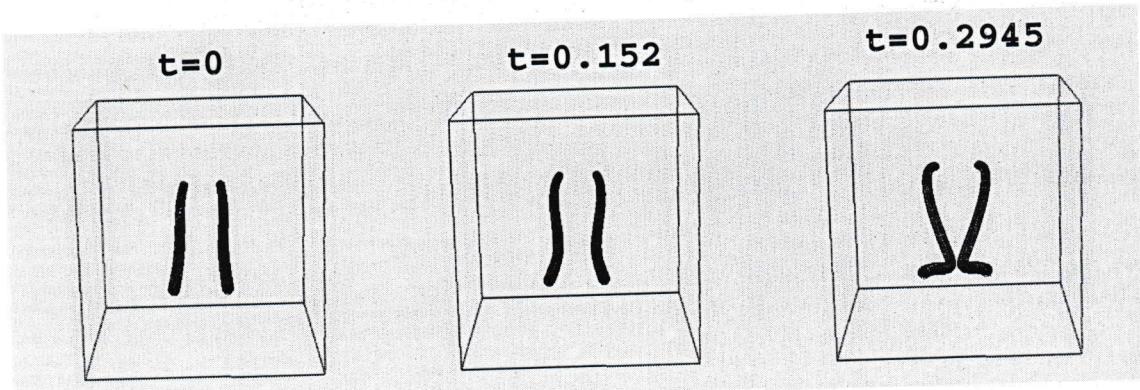
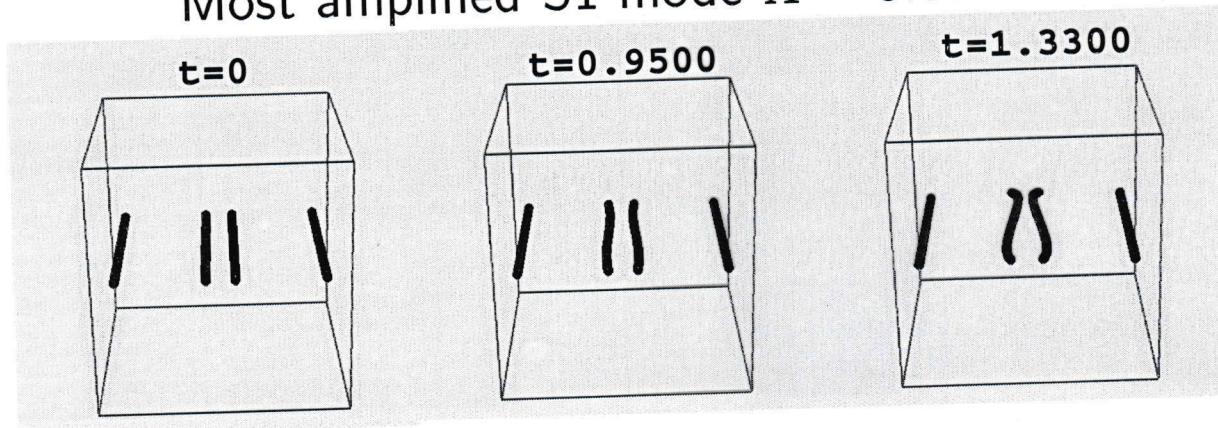


Figure 1: $\Lambda = 10.21$, initial amplitude $\rho = 0.05$, initial thickness $\epsilon = 0.02$ and initial angle $\theta(t = 0) = 47.63(\text{deg})$.

Most amplified S1 mode $\Lambda = 0.8976$.



Long-wave S1 mode $\Lambda = 7.85$.

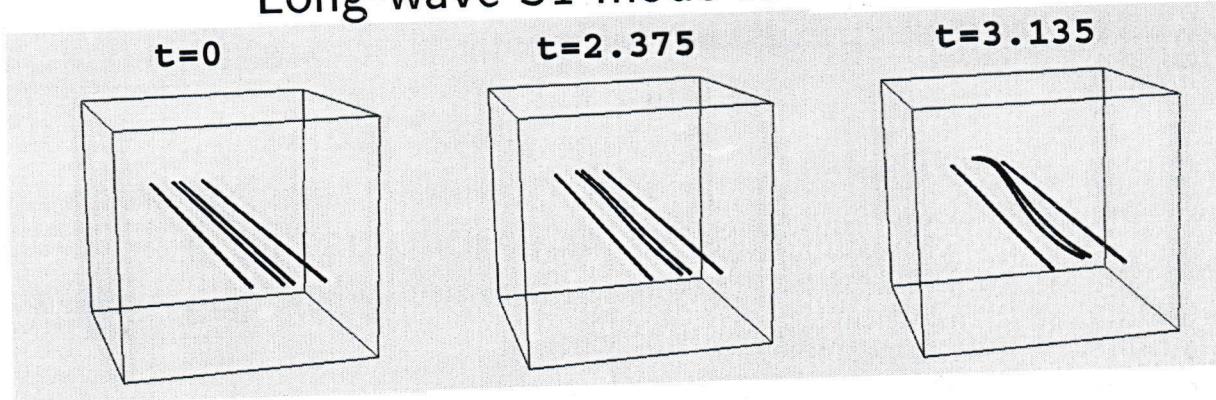


Figure 2: Initial amplitude $\rho_0 = 0.001$ and initial thickness $\epsilon = 0.1$.

Table 1: Four-vortex modes: linear stability (th.) and EZ-vortex (num.) results at $\epsilon = 0.1$.

	Λ	β	θ_o (deg)	θ_i (deg)	$\varrho = \rho_i/\rho_o$	ϱ_o/ϱ_i %
Most amplified S1 mode (th.)*	0.8976	2.91	105.86	131.24	57.4	1.7
// (num.)	//	2.94	111.04	130.20	52.8	1.9
Long-wave S1 mode (th.)*	7.85	1.55	145.45	103.85	9.72	10.3
// (num.)	//	1.56	145.68	103.73	9.80	10.2
Long-wave A mode (th.)*	7.85	1.469	116.90	167.03	9.58	
// (num.)	//	1.511	118.72	166.39	9.73	

*results given by D. Fabre.

Table 2: Four-vortex modes: linear stability (th.) at $\epsilon = 0.02$.

	Λ	β	θ_o (deg)	θ_i (deg)	$\varrho = \rho_i/\rho_o$	ϱ_o/ϱ_i %
Most amplified S1 mode (th.)*	1.2566	3.07	82.81	132.53	48.5	3.0
Long-wave S1 mode (th.)*	7.85	1.62	140.36	104.35	10.00	40
Long-wave A mode (th.)*	7.85	1.40	110.13	167.54	9.35	10.3

*results given by D. Fabre.

Missing Side

Data analysis

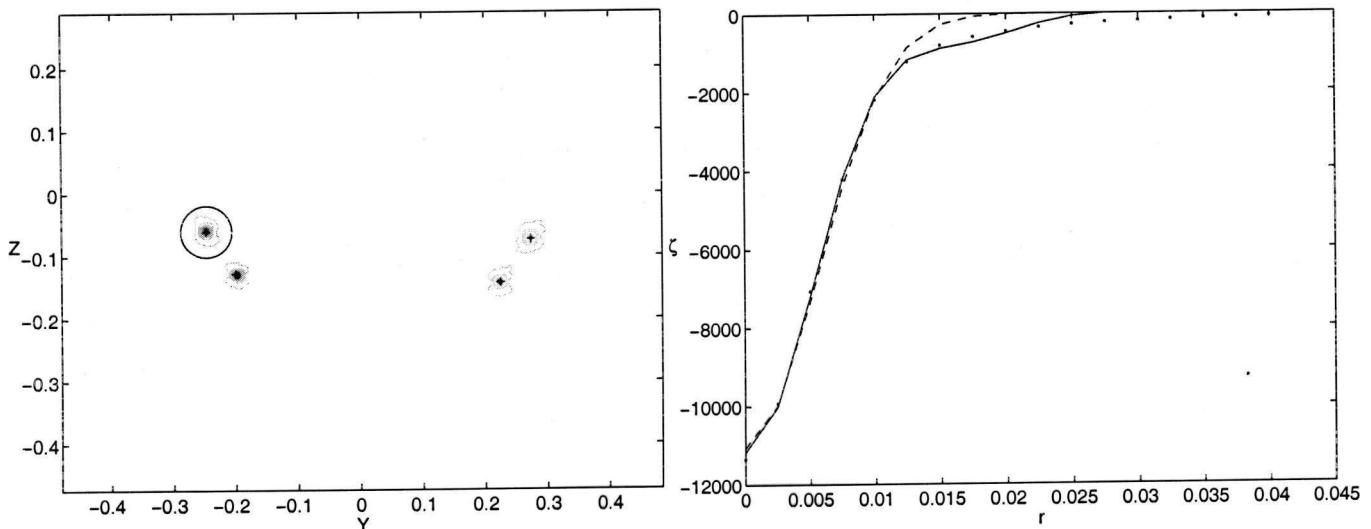


Figure 4: vorticity contour and axisymmetric profil of vorticity

Positions and characteristics of vortices

Wing	x/b	y	z	Γ	δ
left	29.83	-0.281	-0.381	-4.107	0.0237
right	29.83	0.245	-0.344	+4.187	0.0236

Physical parameters of the simulations

$\Gamma (m^2/s)$	$L (m)$	ϵ	$\bar{\nu} (m^2/s)$
± 4.15	0.527	0.0460	0.0077

Numerical computation

Configuration 1

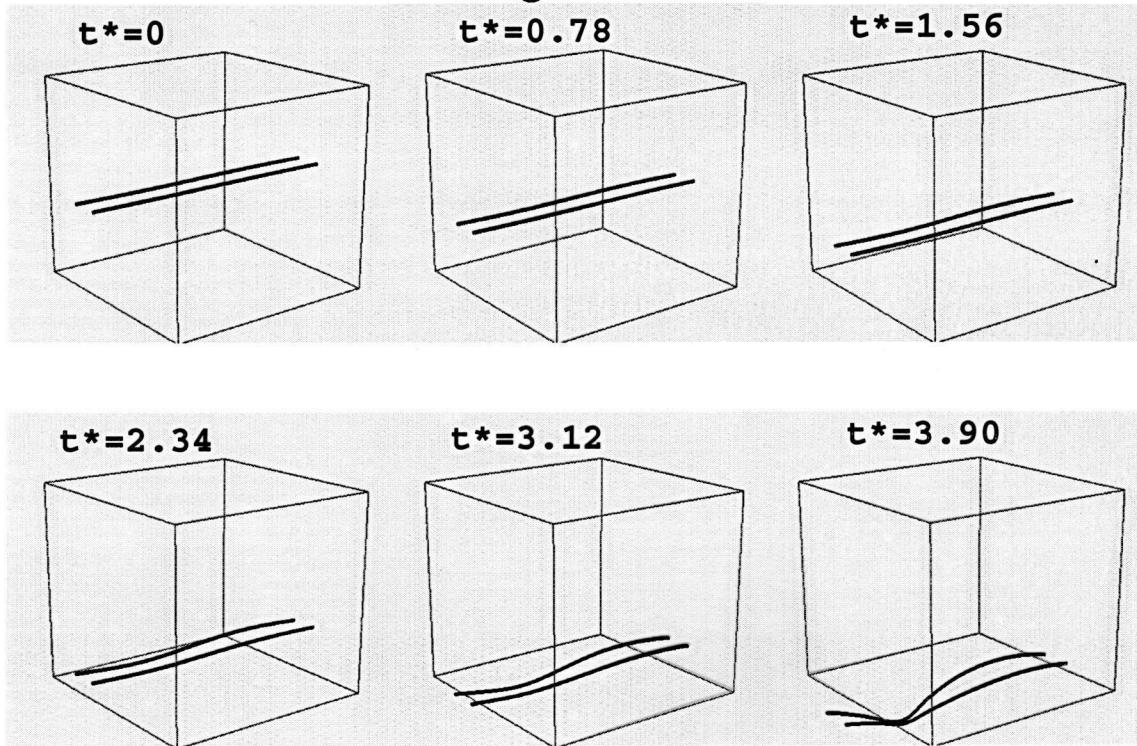


Figure 5: Vortex Filament Simulation of the far wake
for configuration 1.

Fabre, Jacquin and Loof

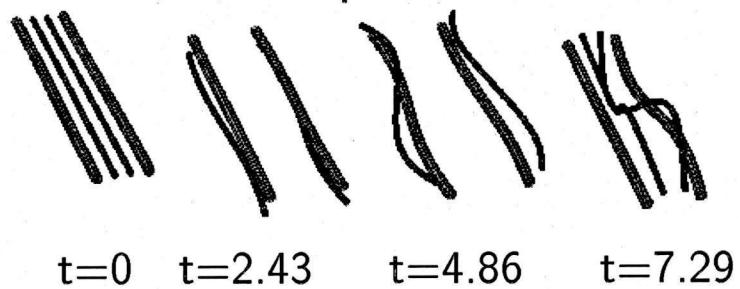
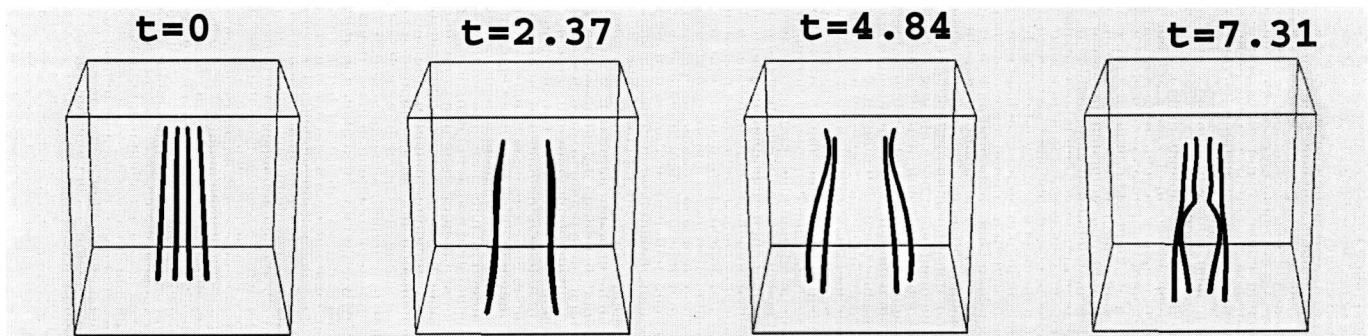
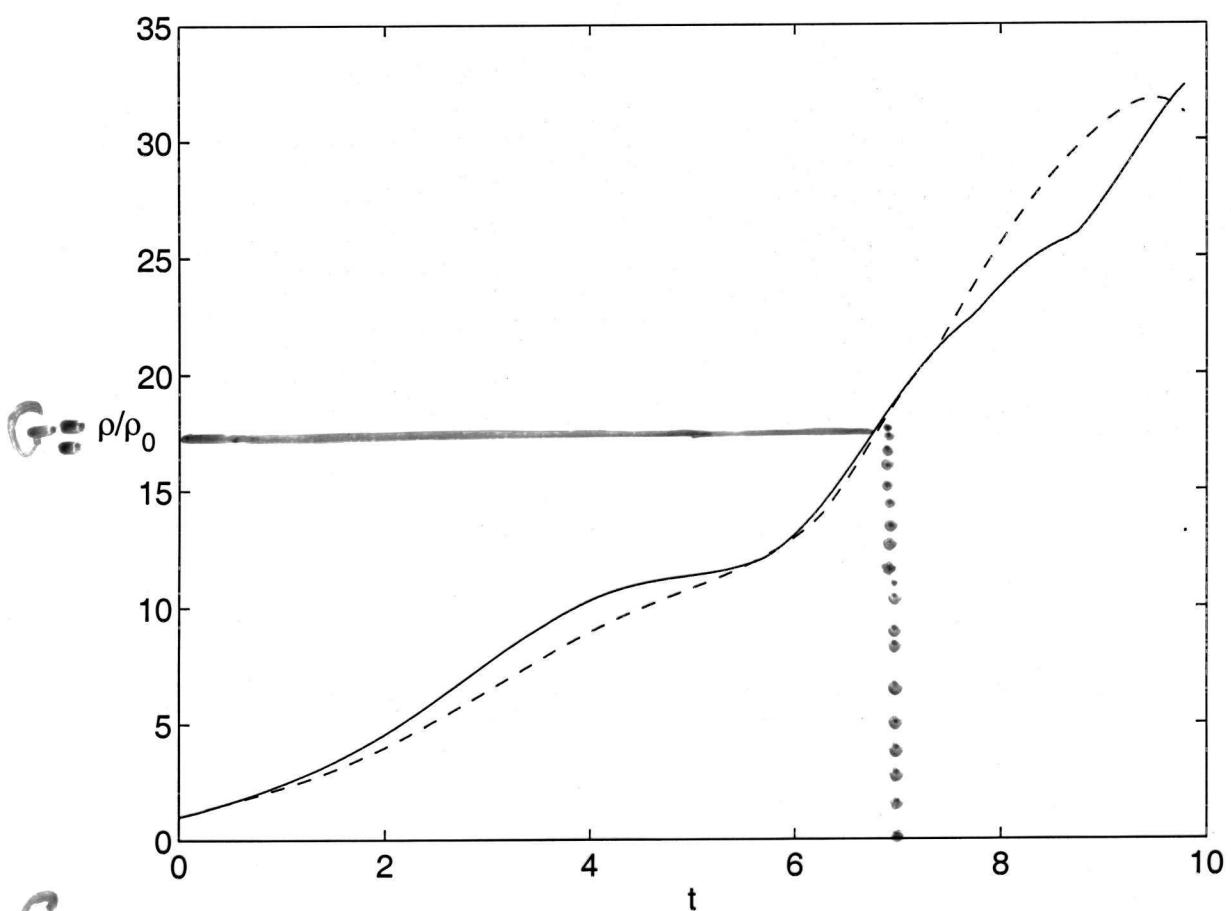
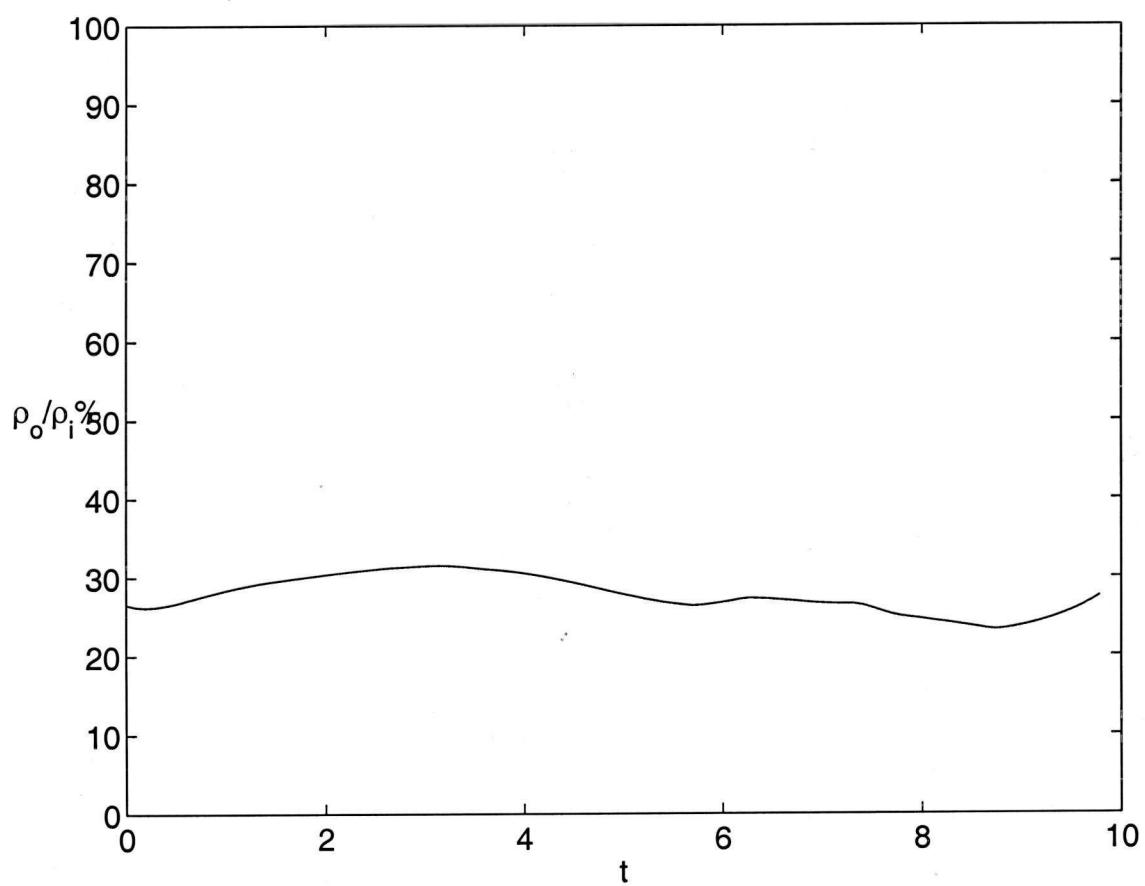


Figure 6: Four vortex wake: $\rho_0 = 0.02$, $\epsilon = 0.1$, $\Lambda = 10.21$, $\Gamma_i/\Gamma_o = -0.3$ and $b_i/b_o = 0.3$





ρ/ρ_0 Faltres Segueia ≈ 14.0



$\rho_0/\rho_i\%$ Faltres Segueia $\approx 24\%$